

SUPPLEMENTARY MATERIA: LOWER BOUND FOR THE EXPECTED SUPREMUM OF FRACTIONAL BROWNIAN MOTION USING COUPLING

KRZYSZTOF BISEWSKI,* Université de Lausanne

Appendix B. Calculations

Derivation of the formula for $\mathcal{E}_{1/2}(T, a)$ in Proposition 2. The derivation in cases $T = \infty$, $a > 0$ and $T \in (0, \infty)$, $a = 0$ is straightforward using Eq. (13) and Eq. (12) respectively. Till this end assume that $a \neq 0$ and $T \in (0, \infty)$. We have

$$\mathbb{P}(\tau(T, a) \in dt) = \left(\frac{e^{-a^2 t/2}}{\sqrt{\pi t}} - \frac{a}{\sqrt{2}} \operatorname{erfc}\left(\frac{a\sqrt{t}}{\sqrt{2}}\right) \right) \left(\frac{e^{-a^2(T-t)/2}}{\sqrt{\pi(T-t)}} + \frac{a}{\sqrt{2}} \operatorname{erfc}\left(-\frac{a\sqrt{T-t}}{\sqrt{2}}\right) \right),$$

see e.g. 2.1.12.4 in [2]. We now have

$$\begin{aligned} \mathcal{E}_{1/2}(T, a) &= \int_0^T t \mathbb{P}(\tau(T, a) \in dt) dt \\ &= T \int_0^1 t \left(\frac{e^{-u^2 t}}{\sqrt{\pi t}} - u \operatorname{erfc}(u\sqrt{t}) \right) \left(\frac{e^{-u^2(1-t)}}{\sqrt{\pi(1-t)}} + u \operatorname{erfc}(-u\sqrt{1-t}) \right) dt, \end{aligned}$$

where we substituted $t = tT$ and put $u := a\sqrt{T/2}$. Now, we have

$$\mathcal{E}_{1/2}(T, a) = T \cdot \left(\frac{e^{-u^2}}{\pi} \cdot J_1 + \frac{u}{\sqrt{\pi}} \cdot (J_2(u) - J_3(u)) - u^2 J_4(u) \right), \quad (38)$$

where

$$\begin{aligned} J_1 &:= \int_0^1 \sqrt{\frac{t}{1-t}} dt = \frac{\pi}{2}, \quad J_2(u) := \int_0^1 \sqrt{t} e^{-u^2 t} \operatorname{erfc}(-u\sqrt{1-t}) dt, \\ J_3(u) &:= \int_0^1 \frac{t}{\sqrt{1-t}} e^{-u^2(1-t)} \operatorname{erfc}(u\sqrt{t}) dt, \quad J_4(u) := \int_0^1 t \operatorname{erfc}(u\sqrt{t}) \operatorname{erfc}(-u\sqrt{1-t}) dt. \end{aligned}$$

Received 25 January 2022; revision received 2 November 2022.

* Postal address: Quartier UNIL-Chamberonne, Bâtiment Extranef, 1015 Lausanne, Switzerland

Using the fact that $\text{erfc}(-z) = 2 - \text{erfc}(z)$ and applying substitution $t = 1 - t$, we obtain

$$J_2(u) = 2 \int_0^1 \sqrt{t} e^{-u^2 t} dt - \int_0^1 \sqrt{1-t} e^{-u^2(1-t)} \text{erfc}(u\sqrt{t}) dt.$$

In the following, let

$$J_5(u) := \int_0^1 \sqrt{t} e^{-u^2 t} dt = \frac{\sqrt{\pi} \text{erf}(u) - 2ue^{-u^2}}{2u^3}, \quad J_6(u) := \int_0^1 \frac{e^{-u^2(1-t)}}{\sqrt{1-t}} \text{erfc}(u\sqrt{t}) dt,$$

where the first integral can be calculated using substitution $t = x^2$. Then

$$\begin{aligned} J_2(u) &= 2J_5(u) - \int_0^1 \frac{1-t}{\sqrt{1-t}} e^{-u^2(1-t)} \text{erfc}(u\sqrt{t}) dt \\ &= 2J_5(u) + J_3(u) - J_6(u). \end{aligned}$$

Now, applying substitution $t = x^2/u^2$ and formula [3, 4.3.20] we obtain

$$J_6(u) = u^{-1} e^{-u^2} \int_0^u \frac{2xe^{x^2}}{\sqrt{u^2 - x^2}} \text{erfc}(ux) dx = \frac{\sqrt{\pi}}{u} \cdot \left(e^{-u^2} - \text{erfc}(u) \right).$$

Further, using integration by parts we obtain

$$\begin{aligned} J_3(u) &= \frac{-2ue^{-u^2(1-t)}\sqrt{1-t} + \sqrt{\pi}(1-2u^2)\text{erf}(u\sqrt{1-t})}{2u^3} \cdot \text{erfc}(u\sqrt{t}) \Big|_0^1 \\ &\quad + \int_0^1 \frac{-2ue^{-u^2(1-t)}\sqrt{1-t} + \sqrt{\pi}(1-2u^2)\text{erf}(u\sqrt{1-t})}{2u^3} \cdot \frac{ue^{-u^2 t}}{\sqrt{\pi t}} dt. \end{aligned}$$

In the following let

$$J_7(u) := \int_0^1 \frac{e^{-u^2 t}}{\sqrt{t}} dt = \frac{\sqrt{\pi} \text{erf}(u)}{u},$$

which can be easily calculated using substitution $t = x^2$. Then

$$J_3(u) = \frac{2ue^{-u^2} - \sqrt{\pi}(1-2u^2)\text{erf}(u)}{2u^3} - \frac{e^{-u^2}}{u\sqrt{\pi}} \cdot J_1 + \left(\frac{1}{2u^2} - 1 \right) (J_7(u) - J_6(u))$$

and therefore

$$\begin{aligned} J_3(u) &= \frac{e^{-u^2}}{2u^3} \left(2u - \sqrt{\pi}(1-u^2) - e^{u^2} \sqrt{\pi}(2u^2-1) \text{erfc}(u) \right) \\ J_2(u) &= \frac{e^{-u^2}}{2u^3} \left(-2u - \sqrt{\pi}(1+u^2) + e^{u^2} \sqrt{\pi}(1+\text{erf}(u)) \right). \end{aligned} \tag{39}$$

In the following, let

$$\begin{aligned} J_8(u) &:= \int_0^1 t \operatorname{erfc}(u\sqrt{t}) dt = \frac{1}{8} \left(4 - \frac{2e^{-u^2}(3+2u^2)}{\sqrt{\pi}u^3} + \left(\frac{3}{u^4} - 4 \right) \operatorname{erf}(u) \right), \\ J_9(u) &:= \int_0^1 t \operatorname{erfc}(u\sqrt{t}) \operatorname{erfc}(-u\sqrt{1-t}) dt, \end{aligned}$$

where the first integral was calculated using integration by parts. Using the identity $\operatorname{erfc}(-z) = 2 - \operatorname{erfc}(z)$ we find that

$$J_4(u) = 2J_8(u) - J_9(u).$$

We will now find the value of the integral $J_9(u)$. See that After applying substitution $t = 1 - t$ we obtain

$$\begin{aligned} 2J_9(u) &= \int_0^1 t \operatorname{erfc}(u\sqrt{t}) \operatorname{erfc}(u\sqrt{1-t}) dt + \int_0^1 (1-t) \operatorname{erfc}(u\sqrt{t}) \operatorname{erfc}(u\sqrt{1-t}) dt \\ &= \int_0^1 \operatorname{erfc}(u\sqrt{t}) \operatorname{erfc}(u\sqrt{1-t}) dt. \end{aligned}$$

After integration by parts we find that

$$\begin{aligned} 2J_9(u) &= \left(\frac{\operatorname{erf}(u\sqrt{t})}{2u^2} + t \operatorname{erfc}(u\sqrt{t}) - \frac{\sqrt{t}e^{-u^2t}}{\sqrt{\pi}u} \right) \cdot \operatorname{erfc}(u\sqrt{1-t}) \Big|_0^1 \\ &\quad - \int_0^1 \left(\frac{\operatorname{erf}(u\sqrt{t})}{2u^2} + t \operatorname{erfc}(u\sqrt{t}) - \frac{\sqrt{t}e^{-u^2t}}{\sqrt{\pi}u} \right) \cdot \frac{ue^{-u^2(1-t)}}{\sqrt{\pi(1-t)}} dt, \end{aligned}$$

therefore

$$2J_9(u) = \left(\frac{\operatorname{erf}(u)}{2u^2} + \operatorname{erfc}(u) - \frac{e^{-u^2}}{\sqrt{\pi}u} \right) - \frac{1}{2\sqrt{\pi}u} (J_7(u) - J_6(u)) - \frac{u}{\sqrt{\pi}} \cdot J_3(u) + \frac{e^{-u^2}}{\pi} \cdot J_1.$$

After simple algebraic manipulations this yields

$$J_4(u) = \frac{3 - e^{-u^2} \left(2u^2 + \frac{6u}{\sqrt{\pi}} \right) + (2u^2 - 3) \operatorname{erfc}(u)}{4u^4}. \quad (40)$$

Finally, after plugging in the expressions for $J_2(u)$, $J_3(u)$ and $J_4(u)$ calculated in (39) and (40) into Eq. (38) we obtain

$$\mathcal{E}_{1/2}(T, a) = \frac{T}{2} \left(1 + \left(\frac{1}{2u^2} - 1 \right) \operatorname{erf}(u) - \frac{e^{-u^2}}{\sqrt{\pi}u} \right),$$

which concludes the proof. \square

Continuation of the proof of Lemma 1. First we show that

$$I_H(t, y) = \frac{t^{H+1/2}}{\sqrt{2\pi}y} \int_0^\infty \int_0^\infty \frac{q^{H-3/2}x^2}{(1+q)^{H+3/2}} \cdot \left[e^{-\frac{1}{2}\left(x-y\sqrt{\frac{q(s)}{t}}\right)^2} - e^{-\frac{1}{2}\left(x+y\sqrt{\frac{q(s)}{t}}\right)^2} \right] dx dq, \quad (41)$$

where $q(s) = s/(1-s)$, which generalizes [1, Proposition 2.1] in case $H = \frac{1}{2}$.

From (41) we have

$$\begin{aligned} I_H(t, y) &= \frac{t^{3/2}}{y} \int_0^t \int_0^\infty \frac{x^2}{s^{3-H}\sqrt{2\pi(t-s)}} e^{-\frac{x^2}{2s} + \frac{y^2}{2t}} \cdot \left[e^{-\frac{(x-y)^2}{2(t-s)}} - e^{-\frac{(x+y)^2}{2(t-s)}} \right] dx ds \\ &= \frac{t^{3/2}}{\sqrt{2\pi}y} \int_0^t \int_0^\infty \frac{x^2}{s^{3-H}\sqrt{t-s}} \cdot \left[e^{-\frac{1}{2}\left(\frac{x}{\sigma_0}-y\mu_0\right)^2} - e^{-\frac{1}{2}\left(\frac{x}{\sigma_0}+y\mu_0\right)^2} \right] dx ds, \end{aligned}$$

where $\mu_0^2 := \frac{s}{(t-s)t}$, and $\sigma_0^2 := \frac{(t-s)s}{t}$. We now apply substitution $x := \sigma_0 x$, then $s := ts$, which gives us

$$I_H(t, y) = \frac{t^{H+1/2}}{\sqrt{2\pi}y} \int_0^1 \int_0^\infty \frac{(1-s)x^2}{s^{3/2-H}} \cdot \left[e^{-\frac{1}{2}\left(x-y\sqrt{\frac{q(s)}{t}}\right)^2} - e^{-\frac{1}{2}\left(x+y\sqrt{\frac{q(s)}{t}}\right)^2} \right] dx ds,$$

which after substitution $q = s/(1-s)$ yields (41). Using [1, Eq. (D.2)] we find that for any $b \in \mathbb{R}$:

$$\int_0^\infty x^2 \left(e^{-\frac{(x-b)^2}{2}} - e^{-\frac{(x+b)^2}{2}} \right) dx = 2be^{-b^2/2} + (1+b^2)\sqrt{2\pi} \cdot \text{erf}(b/\sqrt{2}).$$

Therefore,

$$I_H(t, y) = \frac{t^{H+1/2}}{\sqrt{\pi}y} (J_1(H, u) + J_2(H, u) + J_3(H, u)),$$

where we put $u := \frac{y}{\sqrt{2t}}$ and

$$\begin{aligned} J_1(H, u) &:= \int_0^\infty \frac{q^{H-3/2}}{(1+q)^{H+3/2}} \cdot 2u\sqrt{q}e^{-uq} dq, \\ J_2(H, u) &:= \int_0^\infty \frac{q^{H-3/2}}{(1+q)^{H+3/2}} \cdot \sqrt{\pi} \text{erf}(u\sqrt{q}) dq \\ J_3(H, u) &:= \int_0^\infty \frac{q^{H-3/2}}{(1+q)^{H+3/2}} \cdot 2u^2q^2\sqrt{\pi} \text{erf}(u\sqrt{q}) dq \end{aligned}$$

Using (28) we find that

$$J_1(H, u) = 2u\Gamma(H)U(H, -\frac{1}{2}, u^2).$$

Integration by parts yields

$$J_2(H, u) = \frac{\sqrt{\pi}(2 + 4H + 4q)q^{H-\frac{1}{2}}}{(4H^2 - 1)(1 + q)^{H+1/2}} \cdot \text{erf}(u\sqrt{q}) \Big|_0^\infty - \int_0^\infty \frac{\sqrt{\pi}(2 + 4H + 4q)q^{H-\frac{1}{2}}}{(4H^2 - 1)(1 + q)^{H+1/2}} \cdot \frac{ue^{-u^2q}}{\sqrt{\pi}q} dq$$

and after applying (28) we find that

$$J_2(H, u) = \frac{\sqrt{\pi}}{H^2 - \frac{1}{4}} - \frac{u\Gamma(H)U(H, \frac{1}{2}, u^2)}{H - \frac{1}{2}} - \frac{u\Gamma(H)U(1 + H, \frac{3}{2}, u^2)}{H^2 - \frac{1}{4}}.$$

The integral $J_3(H, u)$ is computed analogously to $J_2(H, u)$ and it equals

$$J_3(u, H) = \frac{2\sqrt{\pi}u^2}{H + \frac{1}{2}} - \frac{2u^3\Gamma(1 + H)U(1 + H, \frac{3}{2}, u^2)}{H + \frac{1}{2}}.$$

Now, after simple algebraic manipulations we obtain

$$J_1(H, u) + J_2(H, u) + J_3(H, u) = \frac{\sqrt{\pi}}{H^2 - \frac{1}{4}} \left(1 - \frac{\Gamma(H)}{\sqrt{\pi}} J_4(H, u) \right) + \frac{2\sqrt{\pi}u^2}{H + \frac{1}{2}},$$

where

$$\begin{aligned} J_4(H, u) &= -2u(H^2 - \frac{1}{4})U(H, -\frac{1}{2}, u^2) + u(H + \frac{1}{2})U(H, \frac{1}{2}, u^2) \\ &\quad + uHU(H + 1, \frac{3}{2}, u^2) + 2u^3(H - \frac{1}{2})HU(H + 1, \frac{3}{2}, u^2). \end{aligned}$$

Applying the relation (31) to the last term above we obtain

$$\begin{aligned} J_4(H, u) &= -2u(H^2 - \frac{1}{4}) \left(U(H, -\frac{1}{2}, u^2) + HU(H + 1, \frac{1}{2}, u^2) \right) + u(H + \frac{1}{2})U(H, \frac{1}{2}, u^2) \\ &\quad + uHU(H + 1, \frac{3}{2}, u^2) + 2u(H - \frac{1}{2})HU(H, \frac{1}{2}, u^2). \end{aligned}$$

Applying the relation (32) to the first and third terms above we obtain

$$J_4(H, u) = uU(H, \frac{3}{2}, u^2) = U(H - \frac{1}{2}, \frac{1}{2}, u^2),$$

where in the last equality we applied the Kummer's transformation (30). This concludes the proof. \square

Lemma 4. *If $H \in (0, \frac{1}{2}) \cup (\frac{1}{2}, 1)$ and $u \in \mathbb{R}$, then*

$$\begin{aligned} &\int_0^\infty U(H - \frac{1}{2}, \frac{1}{2}, z^2) e^{-(z+u)^2} dz \\ &= \frac{\sqrt{\pi}e^{-u^2}}{2\Gamma(H + \frac{1}{2})} + \frac{\sqrt{\pi}|u|^{1-2H}}{2} \left(\frac{\gamma(H + \frac{1}{2}, u^2)}{\Gamma(H + \frac{1}{2})} - \text{sgn}(u) \frac{\gamma(H, u^2)}{\Gamma(H)} \right). \end{aligned}$$

Proof of Lemma 4. Using (29) we find that $U(H - \frac{1}{2}, \frac{1}{2}, z^2) = zU(H, \frac{3}{2}, z^2)$ therefore, using (28) we have

$$\begin{aligned} A(u, H) &:= \int_0^\infty U(H - \frac{1}{2}, \frac{1}{2}, z^2) e^{-(z+u)^2} dz \\ &= \frac{1}{\Gamma(H)} \int_0^\infty t^{H-1} (1+t)^{1/2-H} \int_0^\infty z e^{-(z+u)^2} e^{-z^2 t} dz dt \\ &= \frac{e^{-u^2}}{\Gamma(H)} \int_0^\infty t^{H-1} (1+t)^{1/2-H} \int_0^\infty z e^{-z^2(1+t)-2zu} dz dt \\ &= \frac{e^{-u^2}}{\Gamma(H)} \int_0^\infty t^{H-1} (1+t)^{1/2-H} \cdot \frac{\sqrt{1+t} - \sqrt{\pi} u e^{u^2/(1+t)} \operatorname{erfc}(u/\sqrt{1+t})}{2(1+t)^{3/2}} dt \end{aligned}$$

After substitution $s = \frac{1}{1+t}$ we obtain

$$A(u, H) = \frac{e^{-u^2}}{2\Gamma(H)} (J_1(u, H) - J_2(u, H) + J_3(u, H)),$$

where

$$\begin{aligned} J_1(u, H) &:= \int_0^1 s^{-1/2} (1-s)^{H-1} ds, \\ J_2(u, H) &:= \sqrt{\pi} u \int_0^1 (1-s)^{H-1} e^{u^2 s} ds, \\ J_3(u, H) &:= \sqrt{\pi} u \int_0^1 (1-s)^{H-1} e^{u^2 s} \operatorname{erf}(u\sqrt{s}) ds. \end{aligned}$$

Now, $J_1(u, H)$ can be easily found from the definition of the Beta function, while $J_2(u, H)$ can be found using (27), hence

$$J_1(u, H) = \frac{\sqrt{\pi}\Gamma(H)}{\Gamma(H + \frac{1}{2})}, \quad J_2(u, H) = \frac{\sqrt{\pi}u\Gamma(H)}{\Gamma(H + 1)} {}_1F_1(1, H + 1, u^2).$$

The integral $J_3(u, H)$ is now computed using the error function representation from (37) and the result in (36), i.e.

$$\begin{aligned} J_3(u, H) &= 2u^2 \int_0^1 s^{1/2} (1-s)^{H-1} {}_1F_1(\frac{1}{2}, \frac{3}{2}, u^2 s) ds \\ &= 2u^2 \cdot \frac{\Gamma(\frac{3}{2})\Gamma(H)}{\Gamma(H + \frac{3}{2})} {}_1F_1(1, H + \frac{3}{2}, u^2). \end{aligned}$$

Finally, using (36) we obtain

$$\begin{aligned} A(u, H) &= \frac{\sqrt{\pi}}{2} \left(\frac{e^{-u^2}}{\Gamma(H + \frac{1}{2})} - \frac{ue^{-u^2}}{\Gamma(H+1)} {}_1F_1(1, H+1, u^2) + \frac{u^2 e^{-u^2}}{\Gamma(H + \frac{3}{2})} {}_1F_1(1, H + \frac{3}{2}, u^2) \right) \\ &= \frac{\sqrt{\pi}}{2} \left(\frac{e^{-u^2}}{\Gamma(H + \frac{1}{2})} - \frac{u|u|^{-2H}\gamma(H, u^2)}{\Gamma(H)} + \frac{|u|^{-2H+1}\gamma(H + \frac{1}{2}, u^2)}{\Gamma(H + \frac{1}{2})} \right), \end{aligned}$$

which ends the proof. \square

Proof of Eq. (22). Recall that $H \in (0, \frac{1}{2}) \cup (\frac{1}{2}, 1)$, $a \neq 0$ and $T \in (0, \infty)$. We will first show that

$$\mathcal{J}_H^{(1)}(T, a) := \mathbb{E} \left(\tau^{H-\frac{1}{2}} \left(Y(\tau) + a\tau - \frac{\tau}{Y(\tau)} \right) \right) = 0. \quad (42)$$

The joint density $p(t, y; a, T)$ of the pair $(\tau, Y(\tau))$ is well-known, see (12) in Section 2.2. We have $p(t, y; T, a) = p_1(t, y; T, a)p_2(t; T, a)$, where

$$\begin{aligned} p_1(t, y; T, a) &:= yt^{-3/2} \exp \left\{ -\frac{(y+ta)^2}{2t} \right\} \\ p_2(t; T, a) &:= \frac{1}{\pi\sqrt{T-t}} \cdot \left(e^{-a^2(T-t)/2} + a\sqrt{\frac{\pi(T-t)}{2}} \operatorname{erfc} \left(-a\sqrt{\frac{T-t}{2}} \right) \right), \end{aligned}$$

hence

$$\begin{aligned} \mathcal{J}_H^{(1)}(T, a) &= \int_0^t \int_0^\infty t^{H-1/2} (y + at - \frac{t}{y}) p(t, y; a, T) dy dt \\ &= \int_0^t t^{H-2} p_2(t; T, a) \int_0^\infty (y^2 + aty - t) \exp \left\{ -\frac{(y+at)^2}{2t} \right\} dy dt \\ &= \int_0^t t^{H-1/2} p_2(t; T, a) \int_{a\sqrt{t}}^\infty (z^2 - a\sqrt{t}z - 1) \exp \left\{ -\frac{z^2}{2} \right\} dz dt, \end{aligned}$$

where in the last line we substituted $z := \frac{y+at}{\sqrt{t}}$. It can be seen that the definite integral with respect to z equals to 0 because for every $u \in \mathbb{R}$ we have

$$\begin{aligned} \int_u^\infty z^2 e^{-\frac{z^2}{2}} dz &= ue^{-\frac{u^2}{2}} + \sqrt{\frac{\pi}{2}} \operatorname{erfc} \left(\frac{u}{\sqrt{2}} \right), \quad \int_u^\infty e^{-\frac{z^2}{2}} dz = ue^{-\frac{u^2}{2}}, \\ \int_u^\infty z^2 e^{-\frac{z^2}{2}} dz &= \sqrt{\frac{\pi}{2}} \operatorname{erfc} \left(\frac{u}{\sqrt{2}} \right). \end{aligned}$$

We have established that Eq. (42) holds and therefore it is left to calculate the integral

$$\mathcal{J}_H^{(2)}(T, a) = \frac{\Gamma(H)}{(H + \frac{1}{2})\sqrt{\pi}} \int_0^t t^{H-1} p_2(t; T, a) \int_0^\infty U \left(H - \frac{1}{2}, \frac{1}{2}, \frac{y^2}{2t} \right) \exp \left\{ -\frac{(y+at)^2}{2t} \right\} dy dt.$$

Consider the innermost integral. After substituting $y = z\sqrt{2t}$ we find that

$$\int_0^\infty U\left(H - \frac{1}{2}, \frac{1}{2}, \frac{y^2}{2t}\right) e^{-\frac{(y+at)^2}{2t}} dy = \sqrt{2t} \int_0^\infty U\left(H - \frac{1}{2}, \frac{1}{2}, z^2\right) e^{-(z+a\sqrt{t/2})^2} dz.$$

Using Lemma 4 we obtain

$$\mathcal{J}_H^{(2)}(T, a) = \frac{2^H |a|^{-2H} \Gamma(H)}{\sqrt{2\pi}(H + \frac{1}{2})} \left(J_1(T, a) + J_2(T, a) - J_3(T, a) \right), \quad (43)$$

where

$$\begin{aligned} J_1(T, a) &:= \frac{\sqrt{\pi} 2^{-H} |a|^{2H}}{\Gamma(H + \frac{1}{2})} \int_0^T p_2(t; T, a) t^{H-1/2} e^{-\frac{a^2 t}{2}} dt, \\ J_2(T, a) &:= \frac{\sqrt{\pi} |a|}{\sqrt{2} \Gamma(H + \frac{1}{2})} \int_0^T p_2(t; T, a) \gamma(H + \frac{1}{2}, \frac{a^2 t}{2}) dt, \\ J_3(T, a) &:= \frac{\sqrt{\pi} a}{\sqrt{2} \Gamma(H)} \int_0^T p_2(t; T, a) \gamma(H, \frac{a^2 t}{2}) dt. \end{aligned}$$

It is left to calculate each of the integrals above. Using substitution $t = tT$ and putting $u := a\sqrt{T/2}$ we can find that

$$\begin{aligned} J_1(T, a) &= \frac{|u|^{2H}}{u \Gamma(H + \frac{1}{2})} \int_0^1 \left(\frac{ue^{-u^2(1-t)}}{\sqrt{\pi(1-t)}} + u^2 \operatorname{erfc}(-u\sqrt{1-t}) \right) t^{H-1/2} e^{-u^2 t} dt, \\ J_2(T, a) &= \frac{\operatorname{sgn}(u)}{\Gamma(H + \frac{1}{2})} \int_0^1 \left(\frac{ue^{-u^2(1-t)}}{\sqrt{\pi(1-t)}} + u^2 \operatorname{erfc}(-u\sqrt{1-t}) \right) \gamma(H + \frac{1}{2}, u^2 t) dt, \\ J_3(T, a) &= \frac{1}{\Gamma(H)} \int_0^1 \left(\frac{ue^{-u^2(1-t)}}{\sqrt{\pi(1-t)}} + u^2 \operatorname{erfc}(-u\sqrt{1-t}) \right) \gamma(H, u^2 t) dt. \end{aligned}$$

Let us define:

$$f(t; u) := -\operatorname{erfc}(-u\sqrt{1-t}), \quad g(t; u, H) := -|u|^{-2H-1} \Gamma(H + \frac{1}{2}, u^2 t).$$

Slightly abusing notation, for brevity we write $f(t) := f(t; u)$ and $g(t) := g(t; u, H)$. We then have

$$f'(t) = \frac{ue^{-u^2(1-t)}}{\sqrt{\pi(1-t)}}, \quad g'(t) = t^{H-1/2} e^{-u^2 t},$$

and the quantities $J_i(T, a)$, $i \in \{1, 2, 3\}$ can be expressed as

$$\begin{aligned} J_1(T, a) &= \frac{|u|^{2H}}{u\Gamma(H + \frac{1}{2})} \int_0^1 (f'(t) - u^2 f(t)) g'(t) dt, \\ J_2(T, a) &= \frac{1}{\Gamma(H + \frac{1}{2})} \int_0^1 (f'(t) - u^2 f(t)) \gamma(H + \frac{1}{2}, u^2 t) dt, \\ J_3(T, a) &= \frac{1}{\Gamma(H)} \int_0^1 (f'(t) - u^2 f(t)) \gamma(H, u^2 t) dt. \end{aligned}$$

Before we calculate the values of $J_i(T, a)$ we introduce two useful functions, for

$s > 0, u \in \mathbb{R}$:

$$\begin{aligned} h_1(u, s) &:= \int_0^1 f'(t) \gamma(s, u^2 t) dt = \text{sgn}(u) \cdot \frac{\Gamma(s)\gamma(s + \frac{1}{2}, u^2)}{\Gamma(s + \frac{1}{2})}, \\ h_2(u, s) &:= \int_0^1 (1-t)f'(t) \gamma(s, u^2 t) dt = \text{sgn}(u) \cdot \frac{\Gamma(s)\gamma(s + \frac{3}{2}, u^2)}{2u^2\Gamma(s + \frac{3}{2})}. \end{aligned}$$

The values of these functions were found by applying relation (36) and finding the value of the integral using (33). Now, integration by parts yields

$$\begin{aligned} J_1(T, a) &= \frac{|u|^{2H} e^{-u^2}}{\Gamma(H+1)} + \frac{u|u|^{2H}}{\Gamma(H + \frac{1}{2})} \cdot \left(f(t)g(t) \Big|_0^1 - \int_0^1 f'(t)g(t) dt \right) \\ &= \frac{|u|^{2H} e^{-u^2}}{\Gamma(H+1)} + \frac{\text{sgn}(u)}{\Gamma(H + \frac{1}{2})} \left(\gamma(H + \frac{1}{2}, u^2) + h_1(u, H + \frac{1}{2}) \right) \\ &= \frac{\gamma(H, u^2)}{\Gamma(H)} - \frac{\gamma(H+1, u^2)}{\Gamma(H+1)} + \frac{\text{sgn}(u)}{\Gamma(H + \frac{1}{2})} \left(\gamma(H + \frac{1}{2}, u^2) + h_1(u, H + \frac{1}{2}) \right), \end{aligned}$$

where in the last line we used the recurrence relation for the incomplete Gamma function in Eq. (35). Now, notice that for any $s > 0, u \in \mathbb{R}$:

$$\int \gamma(s, u^2 t) dt = t\gamma(s, u^2 t) + u^{-2}\Gamma(s+1, u^2 t) + C.$$

Therefore, integration by parts yields

$$\begin{aligned} \int_0^1 f(t)\gamma(s, u^2 t) dt &= -\gamma(s, u^2) - \int_0^1 t f'(t) \gamma(s, u^2 t) dt + u^{-2} \int_0^1 f'(t) \gamma(s+1, u^2 t) dt \\ &= -\gamma(s, u^2) + h_2(u, s) - h_1(u, s) + u^{-2}h_1(u, s+1). \end{aligned}$$

Finally, this gives us

$$J_2(T, a) = \frac{\operatorname{sgn}(u)}{\Gamma(H + \frac{1}{2})} \left((1+u^2)h_1(u, H + \frac{1}{2}) + u^2\gamma(H + \frac{1}{2}, u^2) - u^2h_2(u, H + \frac{1}{2}) \right. \\ \left. - \gamma(H + \frac{3}{2}) - h_1(u, H + \frac{3}{2}) \right).$$

Using analogous methods, we find that

$$J_3(T, a) = \frac{1}{\Gamma(H)} \left((1+u^2)h_1(u, H) + u^2\gamma(H, u^2) - u^2h_2(u, H) - \gamma(H+1, u^2) - h_1(u, H+1) \right).$$

Through straightforward algebraic manipulations and applying simple recurrence relation (35), we finally obtain

$$J_1(T, a) + J_2(T, a) - J_3(T, a) = \frac{\gamma(H, u^2)}{\Gamma(H)},$$

which concludes the proof. \square

References

- [1] BISEWSKI, K., DĘBICKI, K. AND ROLSKI, T. (2022). Derivatives of sup-functionals of fractional Brownian motion evaluated at $H=1/2$. *Electron. J. Probab.* **27**, 1–35.
- [2] BORODIN, A. N. AND SALMINEN, P. (2002). *Handbook of Brownian motion—facts and formulae* second ed. Probability and its Applications. Birkhäuser Verlag, Basel.
- [3] NG, E. W. AND GELLER, M. (1969). A table of integrals of the error functions. *J. Res. Natl. Bur. Standards Sect. B* **73B**, 1–20.