

# First-order velocity field

## Appendix A. Series representation of the velocity field

The first-order velocity field  $\mathbf{u}^{(1)}$  can be calculated according to the method by Brenner (1964). However, as pointed out by Dörr & Hardt (2015), corrections to Brenner's method are necessary. Here, we provide all of the expressions required to explicitly compute the velocity field  $\mathbf{u}^{(1)}$ , which is of the form (Brenner 1964)

$$\mathbf{u}^{(1)} = \sum_{k=0}^{\infty} \mathbf{u}_k^{(1)}. \quad (\text{A } 1)$$

To facilitate the explicit evaluation of the velocity field and to simultaneously ensure a high degree of accuracy, we truncate the series after 20 summation terms and arrive at

$$\mathbf{u}^{(1)} \approx \sum_{k=0}^{20} \mathbf{u}_k^{(1)}. \quad (\text{A } 2)$$

The velocity fields  $\mathbf{u}_k^{(1)}$  can be expressed by

$$\begin{aligned} \mathbf{u}_k^{(1)} = \sum_{n=1}^{\infty} & \left[ \nabla \times \left( \mathbf{r}_k \chi_{-n-1}^{(1)} \right) + \nabla \left( {}_k\phi_{-n-1}^{(1)} \right) - \frac{n-2}{2n(2n-1)\mu_1} r^2 \nabla \left( {}_k p_{-n-1}^{(1)} \right) \right. \\ & \left. + \mathbf{r} \frac{n+1}{n(2n-1)\mu_1} {}_k p_{-n-1}^{(1)} \right], \end{aligned} \quad (\text{A } 3)$$

where

$${}_k p_{-n-1}^{(1)} = \frac{(2n-1)\mu_1}{(n+1)a} \left( \frac{a}{r} \right)^{n+1} \left[ (n+2) {}_k X_n^{(1)} + {}_k Y_n^{(1)} \right], \quad (\text{A } 4)$$

$${}_k \phi_{-n-1}^{(1)} = \frac{a}{2(n+1)} \left( \frac{a}{r} \right)^{n+1} \left( n {}_k X_n^{(1)} + {}_k Y_n^{(1)} \right), \text{ and} \quad (\text{A } 5)$$

$${}_k \chi_{-n-1}^{(1)} = \frac{1}{n(n+1)} \left( \frac{a}{r} \right)^{n+1} {}_k Z_n^{(1)} \quad (\text{A } 6)$$

(Brenner 1964). Through accommodating the velocity field (A 3) to the boundary condition at the particle surface, the functions  ${}_k X_n^{(1)}$ ,  ${}_k Y_n^{(1)}$  and  ${}_k Z_n^{(1)}$  can be specified. One has

$${}_k X_n^{(1)} = 0 \quad \text{for } n \leq 1, \quad (\text{A } 7)$$

$${}_k Y_n^{(1)} = \begin{cases} \frac{3}{2} \sum_{m=0}^k \frac{c_{km} N_{km} \cos(m\varphi)}{2k+1} (k-1)(k+m) P_{k-1}^m(\cos\theta) & \text{for } n = k-1 \\ -\frac{3}{2} \sum_{m=0}^k \frac{c_{km} N_{km} \cos(m\varphi)}{2k+1} (2+k)(k-m+1) P_{k+1}^m(\cos\theta) & \text{for } n = k+1 \\ 0 & \text{for all other } n, \end{cases} \quad (\text{A } 8)$$

and

$${}_k Z_n^{(1)} = \begin{cases} \frac{3}{2} U \sum_{m=0}^k m c_{km} N_{km} P_k^m(\cos\theta) \sin(m\varphi) & \text{for } n = k \\ 0 & \text{for } n \neq k. \end{cases} \quad (\text{A } 9)$$

In equations (A 7)–(A 9), the  $P_k^m$  are associated Legendre polynomials. Expression (A 8), developed by Dörr & Hardt (2015), differs from the corresponding result by Brenner (1964), while equations (A 7) and (A 9) have been adopted from Brenner (1964) without modification. The expansion coefficients  $c_{km}$  and  $N_{km}$  are given by

$$N_{km} = \begin{cases} \sqrt{2} \sqrt{\frac{2k+1}{4\pi}} \frac{(k-m)!}{(k+m)!} & \text{if } m > 0 \\ \sqrt{\frac{2k+1}{4\pi}} & \text{if } m = 0 \end{cases} \quad (\text{A } 10)$$

and

$$c_{km} = \int_0^\pi \int_0^{2\pi} \sin \theta |\cos \varphi| N_{km} P_k^m(\cos \theta) \cos(m\varphi) \sin \theta d\theta d\varphi, \quad (\text{A } 11)$$

respectively. Equation (A 11) contains information about the particle geometry via the term  $\sin \theta |\cos \varphi|$ , which is equal to the first-order term  $\phi^{(1)}$  in the expansion of the particle shape (equation (2.5) in the paper). The velocity field  $\mathbf{u}^{(1)}$  according to equation (A 1) is thus fully determined. For the sake of convenience, we explicitly display the coefficients  $c_{km}$  for  $k \leq 20$  in appendix B. Due to the symmetry of the particle shape with respect to the plane  $x = 0$ , coefficients with  $m < 0$  and/or  $m$  odd vanish.

## Appendix B. Coefficients $c_{km}$

$k=0: m=0$

$$\sqrt{\pi}$$

$k=2: m=0, 2$

$$-\frac{\sqrt{5\pi}}{8}, \frac{\sqrt{15\pi}}{8}$$

$k=4: m=0, 2, 4$

$$-\frac{3\sqrt{\pi}}{64}, \frac{\sqrt{5\pi}}{32}, -\frac{\sqrt{35\pi}}{64}$$

$k=6: m=0, 2, 4, 6$

$$-\frac{5\sqrt{13\pi}}{1024}, \frac{\sqrt{\frac{1365\pi}{2}}}{1024}, -\frac{3\sqrt{91\pi}}{1024}, \frac{\sqrt{\frac{3003\pi}{2}}}{1024}$$

$k=8: m=0, 2, 4, 6, 8$

$$-\frac{35\sqrt{17\pi}}{16384}, \frac{3\sqrt{\frac{595\pi}{2}}}{4096}, -\frac{3\sqrt{1309\pi}}{8192}, \frac{\sqrt{\frac{7293\pi}{2}}}{4096}, -\frac{3\sqrt{12155\pi}}{16384}$$

$k=10: m=0, 2, \dots, 10$

$$-\frac{147\sqrt{21\pi}}{131072}, \frac{49\sqrt{385\pi}}{131072}, -\frac{7\sqrt{5005\pi}}{65536}, \frac{21\sqrt{\frac{5005\pi}{2}}}{131072}, -\frac{7\sqrt{\frac{85085\pi}{3}}}{131072}, \frac{7\sqrt{\frac{323323\pi}{6}}}{131072}$$

$k=12: m=0, 2, \dots, 12$

$$-\frac{3465\sqrt{\pi}}{1048576}, \frac{45\sqrt{3003\pi}}{524288}, -\frac{225\sqrt{\frac{1001\pi}{2}}}{1048576}, \frac{75\sqrt{\frac{2431\pi}{2}}}{524288}, -\frac{15\sqrt{138567\pi}}{1048576}, \frac{15\sqrt{\frac{88179\pi}{2}}}{524288}, -\frac{15\sqrt{\frac{676039\pi}{2}}}{1048576}$$

$k=14: m=0, 2, \dots, 14$

$$-\frac{14157\sqrt{29\pi}}{33554432}, \frac{1089\sqrt{39585\pi}}{67108864}, -\frac{99\sqrt{\frac{2467465\pi}{2}}}{33554432}, \frac{33\sqrt{46881835\pi}}{67108864}, -\frac{33\sqrt{12785955\pi}}{33554432}, \frac{33\sqrt{58815393\pi}}{67108864}, \\ -\frac{165\sqrt{\frac{1508087\pi}{2}}}{33554432}, \frac{495\sqrt{646323\pi}}{67108864}$$

$k=16$ :  $m=0,2,\dots,16$

$$\begin{aligned}
 & -\frac{306735\sqrt{33\pi}}{1073741824}, \frac{20449\sqrt{935\pi}}{268435456}, -\frac{1573\sqrt{\frac{323323\pi}{2}}}{268435456}, \frac{3003\sqrt{46189\pi}}{268435456}, \\
 & -\frac{1001\sqrt{\frac{5311735\pi}{3}}}{536870912}, \frac{715\sqrt{\frac{2860165\pi}{3}}}{268435456}, -\frac{2145\sqrt{\frac{245157\pi}{2}}}{268435456}, \frac{143\sqrt{35547765\pi}}{268435456}, -\frac{143\sqrt{1101980715\pi}}{1073741824}
 \end{aligned}$$

$k=18$ :  $m=0,2,\dots,18$

$$\begin{aligned}
 & -\frac{1738165\sqrt{37\pi}}{8589934592}, \frac{61347\sqrt{59755\pi}}{8589934592}, -\frac{5577\sqrt{\frac{920227\pi}{2}}}{2147483648}, \frac{1001\sqrt{\frac{117920517\pi}{2}}}{4294967296}, -\frac{15015\sqrt{274873\pi}}{4294967296}, \\
 & \frac{2145\sqrt{\frac{28861665\pi}{2}}}{4294967296}, -\frac{2145\sqrt{\frac{7971317\pi}{2}}}{2147483648}, \frac{429\sqrt{\frac{3706662405\pi}{2}}}{8589934592}, -\frac{429\sqrt{2398428615\pi}}{8589934592}, \frac{715\sqrt{\frac{3357800061\pi}{2}}}{8589934592}
 \end{aligned}$$

$k=20$ :  $m=0,2,\dots,20$

$$\begin{aligned}
 & -\frac{10207769\sqrt{41\pi}}{68719476736}, \frac{48841\sqrt{899745\pi}}{34359738368}, -\frac{8619\sqrt{117266765\pi}}{68719476736}, \frac{1105\sqrt{\frac{914680767\pi}{2}}}{17179869184}, -\frac{9945\sqrt{23453353\pi}}{34359738368}, \\
 & \frac{1989\sqrt{\frac{309157835\pi}{2}}}{17179869184}, -\frac{3315\sqrt{\frac{1916778577\pi}{2}}}{68719476736}, \frac{3315\sqrt{\frac{531543639\pi}{2}}}{34359738368}, -\frac{3315\sqrt{1240268491\pi}}{68719476736}, \frac{1105\sqrt{\frac{7245779079\pi}{2}}}{34359738368}, \\
 & -\frac{663\sqrt{\frac{156991880045\pi}{2}}}{68719476736}
 \end{aligned}$$

## REFERENCES

- BRENNER, H. 1964 The Stokes resistance of a slightly deformed sphere. *Chem. Eng. Sci.* **19**, 519–539.
- DÖRR, A. & HARDT, S. 2015 Driven particles at fluid interfaces acting as capillary dipoles. *J. Fluid Mech.* **770**, 5–26.